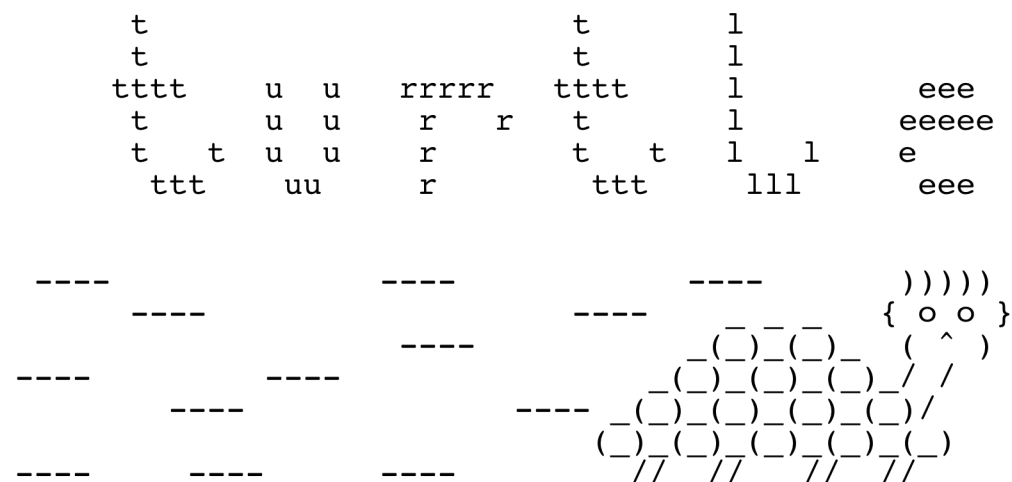


Valence Bond

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Contents

- What is VB / Spinfuctions/Bonds
- Orbital Optimisation / Brillouin Theorem
- Energies and Gradients - Matrix elements
- Applications/Illustrations
 - Is the BSSE problem solved by using Non-orthogonal Orbitals ?
 - What is a π system in a bent benzene ?
 - Resonance and aromaticity
 - Local Resonance
 - Do VB structures mean something ?
 - What about O₂

VB techniques are like Orthogonal methods

What Universe are we in:

- Time independent
- Electronic Structure
- Born Oppenheimer
- Only scalar relativity
- The Answer is 42
- The Question ??

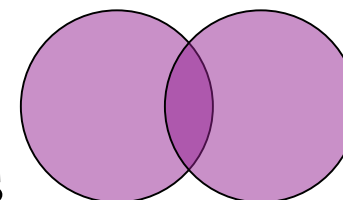
$$H = \sum_i h(i) + \sum_{i < j} \frac{1}{r_{ij}} + N.R. + \sim$$

1-electron

2-electron

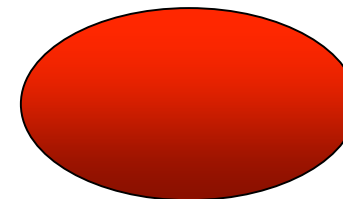
VB

Starting from : Atoms
Bonding : Spin-pairing
Orbitals : Non-Orthogonal
(historically) fixed
Simplest : 1 structure, more determinants
In general : Multi-structure
Extend : Add Ionic.. structures

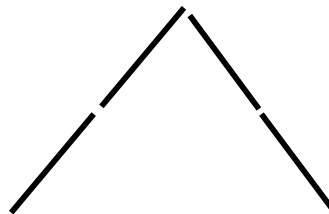
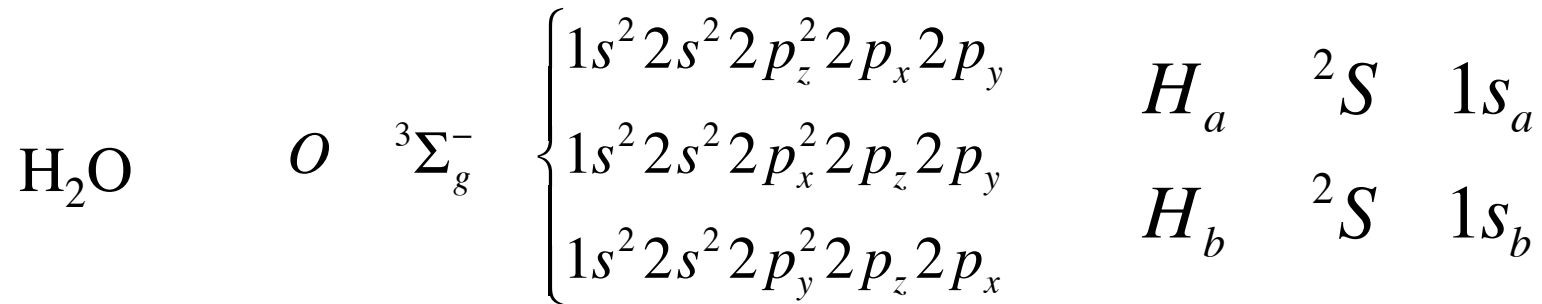


MO

Starting from : Molecules
Bonding : Sharing Orbitals
Orbitals : Orthogonal
Optimised
In general : 1 configuration, 1 determinant
Extend : Add configurations (CI)



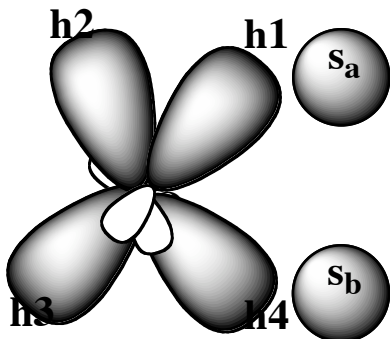
VB-atoms



and



or



$$\Psi = 1s^2 h_2^2 h_3^2 h_1 s_a h_4 s_b$$



SPIN COUPLING

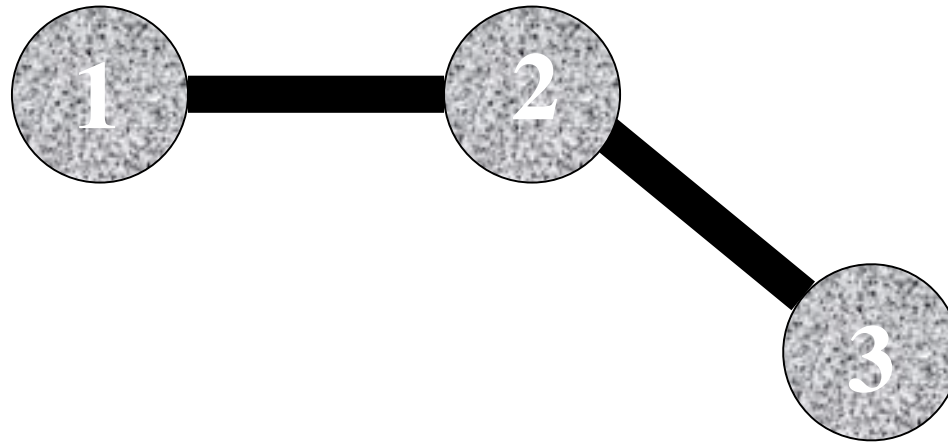
Perfect pairing - two orbitals are coupled to a singlet (bond)



$$\Psi = N(|\varphi_1\bar{\varphi}_2| - |\bar{\varphi}_1\varphi_2|)$$

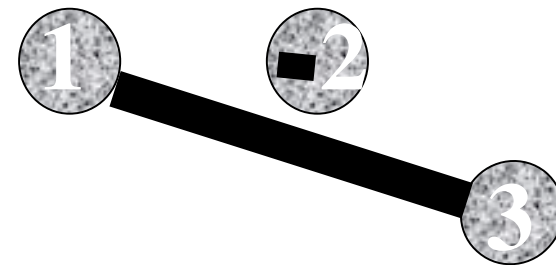
$$\Psi = (\widehat{\varphi_1\varphi_2})$$

Two Bonds



$$\Psi = N \mathbb{A}(|\varphi_1 \bar{\varphi}_{2a}| - |\bar{\varphi}_1 \varphi_{2a}|)(|\varphi_{2b} \bar{\varphi}_3| - |\bar{\varphi}_{2b} \varphi_3|) = (\widehat{\varphi_1 \varphi_{2a}})(\widehat{\varphi_{2b} \varphi_3})$$

$$\Psi = (\widehat{\varphi_1 \varphi_3})(\widehat{\varphi_{2a} \varphi_{2b}})$$

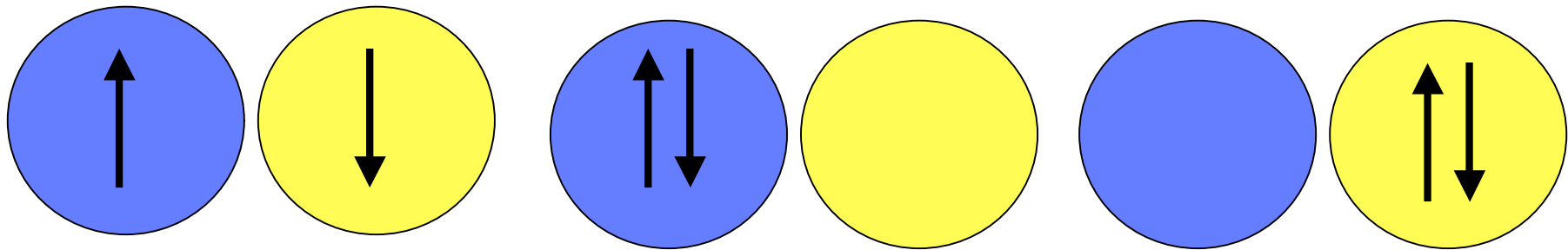


VB-2

H₂ vergeet spin

$$\Psi(1,2) = c_1 \underbrace{(s_a(1)s_b(2) + s_a(2)s_b(1))}_{\text{Covalent}}$$

$$+ c_2 \underbrace{(s_a(1)s_a(2))}_{\text{ionisch}} + c_3 \underbrace{(s_b(1)s_b(2))}_{\text{ionisch}}$$



Covalent

ionisch

ionisch

H₂ VB possibilities

$$\Psi_{\text{singlet}} = \frac{1}{2}(1s_A 1s_B + 1s_B 1s_A)(\alpha\beta - \beta\alpha) \quad m_s = 0, s = 0 \quad (\text{singlet } \uparrow\downarrow)$$
$$\Psi_{\text{triplet}} = \begin{cases} \frac{1}{2}(1s_A 1s_B - 1s_B 1s_A)(\alpha\beta + \beta\alpha) & m_s = 0, s = 1 \quad (\text{triplet } \uparrow\uparrow) \\ \frac{1}{\sqrt{2}}(1s_A 1s_B - 1s_B 1s_A)(\alpha\alpha) & m_s = 1, s = 1 \quad (\text{triplet } \uparrow\uparrow) \\ \frac{1}{\sqrt{2}}(1s_A 1s_B - 1s_B 1s_A)(\beta\beta) & m_s = -1, s = 1 \quad (\text{triplet } \uparrow\uparrow) \end{cases}$$

Volgorde elektronen - 1..2 1..2

Kunt $\langle E \rangle = \frac{\langle \Psi | H | \Psi \rangle}{\langle \Psi | \Psi \rangle}$ berekenen; Grondtoestand is natuurlijk singlet

De Hamiltoniaan bevat (bij ons) enkel ruimte afhankelijke termen. Dus wordt de energie door het ruimte deel bepaald.

VB energy H₂ - 3

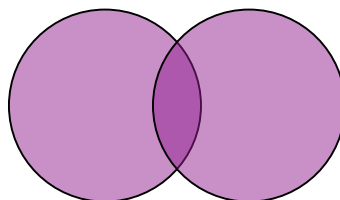
$$\Psi_{\text{singlet}} : \frac{1}{2}(1s_A 1s_B + 1s_B 1s_A) \quad (\text{singlet } \uparrow\downarrow) \quad E \approx E_1 + J + K$$
$$\Psi_{\text{triplet}} : \begin{cases} \frac{1}{2}(1s_A 1s_B - 1s_B 1s_A) & (\text{triplet } \uparrow\uparrow) \\ \frac{1}{\sqrt{2}}(1s_A 1s_B - 1s_B 1s_A) & (\text{triplet } \uparrow\uparrow) \quad E \approx E_1 + J - K \\ \frac{1}{\sqrt{2}}(1s_A 1s_B - 1s_B 1s_A) & (\text{triplet } \uparrow\uparrow) \end{cases}$$

gives bonding

J en K : Bevoordelen triplet (denk He* / Regel van Hund)
VB attractief door 1-elektron term (niet orthogonaal)
But J,K,E are different

Can we use **orthogonal** Orbitals in VB ??

H₂(STO6G)



$$\Psi = \frac{1}{\sqrt{2}} [|a\bar{b}| + |b\bar{a}|]$$

fixed orbitals

optimised orbitals

VB: Heitler-London :	0.70 Å	120.6 Kcal/mol
VB: Coulson-Fischer: (full CI)	0.73 Å	128.0 Kcal/mol
Hartree-Fock :	0.71 Å	115.6 Kcal/mol
Orthogonal VB :	>5.9 Å	0 Kcal/mol
Experimental (WWW):	0.74 Å	104 Kcal/mol

NO

Atomic Population Analysis (cf. McWeeny (1953))

H2 (STO6G/0.70 Å)

	Orb-overlap	q_a	q_{ab}	Energy
VB: H-L	0.69	0.68	0.32	-1.134187
VB: C-F	0.81	0.62	0.38	-1.144979
HF :	0.0	0.60	0.40	-1.126100
Orthog VB :	0.0	1.89	-0.89	-0.463516
VB:C-F (+p)	0.81	0.61	0.39	-1.148056

VB-hierarchy

$$\Psi = \sum C_i \Phi_i \quad \text{L} \quad \text{Wavefunction}$$

$$(H - ES)C = 0 \text{ or fixed}$$

$$\Phi = \sum s_d \Delta_d \quad \text{L} \quad \text{Structure}$$

$$\text{Spin-projection}$$

$$\Delta = |\psi_1 \psi_2 \dots \psi_n| \quad \text{L} \quad \text{Determinant}$$

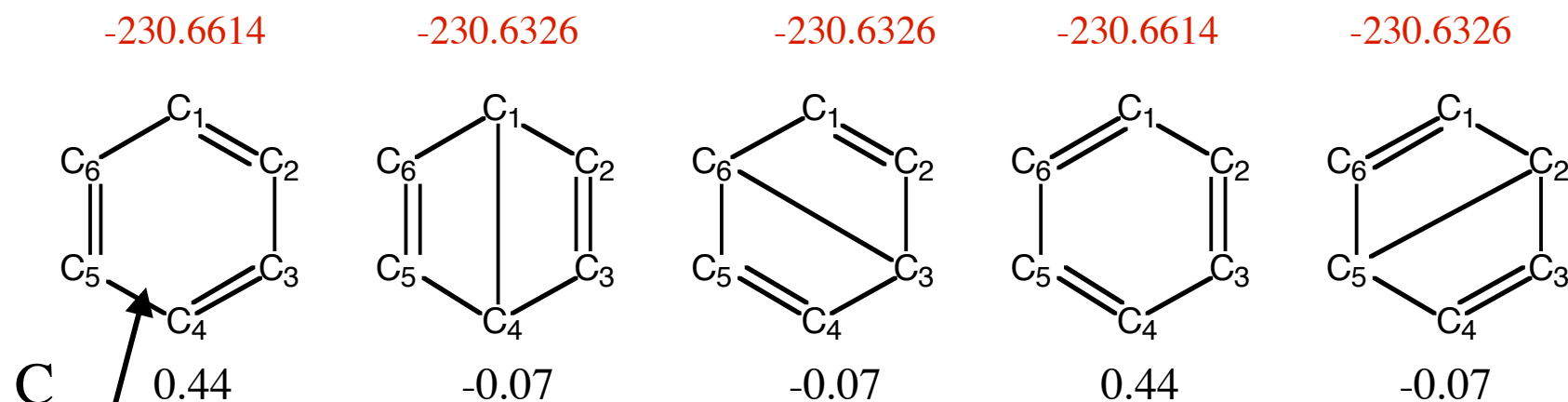
$$\psi = \sum c_i \phi_i \quad \text{L} \quad \text{Orbital}$$

$$\text{Optimised or fixed}$$

basis-function

Spin-Functions(1)

Rumer functions : Different orbitals spin-paired
Benzene (6-31G)



$$E = -230.6938$$

$$\Psi_1 = N(|\varphi_1\bar{\varphi}_2| - |\bar{\varphi}_1\varphi_2|)(|\varphi_3\bar{\varphi}_4| - |\bar{\varphi}_3\varphi_4|)(|\varphi_5\bar{\varphi}_6| - |\bar{\varphi}_5\varphi_6|)$$

Spin-functions non-orthogonal

$$\langle \psi_1 | \psi_2 \rangle = 0.52$$

$$\langle \Psi_1 | \Psi_4 \rangle_{spin} = 0.25$$

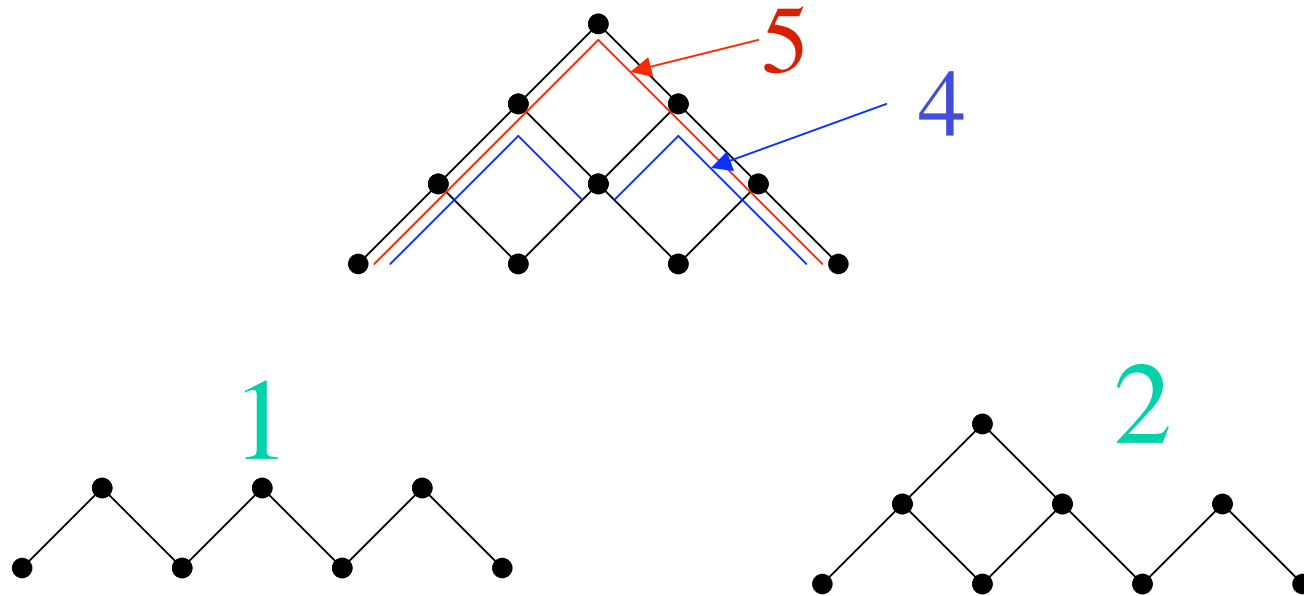
$$\langle \Psi_1 | \Psi_4 \rangle = 0.66$$

$$\langle \Psi_1 | \Psi_3 \rangle_{spin} = 0.5$$

$$\langle \Psi_1 | \Psi_3 \rangle = 0.82$$

Spin-Functions(2)

Branching Diagrams : Orbitals in standard order
Different spin-functions



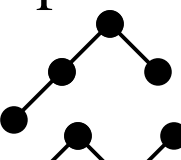
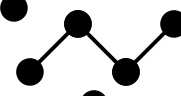
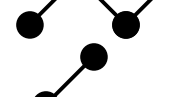

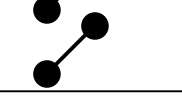
Spin-functions orthogonal;

Obtain by e.g Schmidt orthogonalising Rumer functions

Spin-Functions(3)

⊗

Branching Diagrams are often considered less informative
 OH ($^2\Pi$) (JHvL, G.G.Balint-Kurti (1980) DZP, Atomic Opt

#	Atomic states	Configuration	spin	Coefficient
1	O(3P) H(2S)	$1s^2 2s^2 y^2 xzh$		-0.235
2	O(1D) H(2S)	$1s^2 2s^2 y^2 xyz$		0.191
3	O $^+(^2D^0)$ H $^-(^1S)$	$1s^2 2s^2 (z^2 - y^2) xh^2$		0.053
4	O $^+(^2P^0)$ H $^-(^1S)$	$1s^2 2s^2 (z^2 + y^2) xh^2$		-0.028
5	O $^-(^2P^0)$ H $^+$	$1s^2 2s^2 z^2 y^2 x$		0.327

- Oxygen 1D mixes in
- O $^-$ ionic structure important due to atomic orbitals
- O(z) => hybrid on bond formation
- Spin-function may be transformed from one basis to another

VB=MCSCF

$$\Psi = N(c_1|\varphi_3\bar{\varphi}_3| - c_2|\bar{\varphi}_4\varphi_4|) =$$

$$\text{with } \left\{ \begin{array}{l} \varphi_1 = \frac{(\sqrt{c_1}\varphi_3 + \sqrt{c_2}\varphi_4)}{(c_1 + c_2)} \\ \varphi_2 = \frac{(\sqrt{c_1}\varphi_3 - \sqrt{c_2}\varphi_4)}{(c_1 + c_2)} \end{array} \right\}$$

$$= N(|\varphi_1\bar{\varphi}_2| - |\bar{\varphi}_1\varphi_2|) = (\varphi_1\varphi_2)$$

$$\langle \varphi_1 | \varphi_2 \rangle = \frac{c_1 - c_2}{c_1 + c_2} \quad \langle \varphi_3 | \varphi_4 \rangle = 0$$

$$\widehat{(\phi_1\phi_2)}\widehat{(\phi_3\phi_4)} \overset{PP}{=} (2 - MCSCF) * (2 - MCSCF)$$

Only if $(\phi_1\phi_2)$ orthogonal to $(\phi_3\phi_4)$: **Strong Orthogonality**

GVB : W.A. Goddard III et al (1972 - ..)

Original version: general

Popular version : GVB-PP

strong orthogonality

perfect pairing

uses Fock-operators (N^4)

size extensive

1999 : GVB-RP

Spincoupled / Spincoupled VB

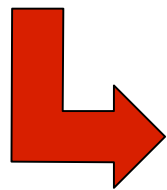
J. Gerratt[†], M. Raimondi, D.L. Cooper - “The gang of Three”

Spincoupled : 1 configuration all spincouplings, optimised

$$\Psi = \sum_k^{f_S^N} c_{Sk} (N!)^{1/2} A \left\{ \psi_1 \psi_2 \psi_3 \dots \psi_N \Theta_{S,M;k}^N \right\}$$

This function can be obtained by restricting a CASSCF function : **CASVB**

Each ψ_i satisfies her own $F_i^{eff} \psi_i = \epsilon_i \psi_i$



“stack” of virtual orbitals ($\psi_i^1, \psi_i^2, \dots$)

Generate excited structures by applying excitations :

$C_{i \rightarrow i1}$ or $C_{i \rightarrow i2}, \dots$ (singles), $C_{i \rightarrow i1, j \rightarrow ij1}$ (double)



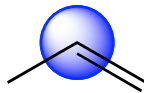
Spincoupled-VB

Orbital models

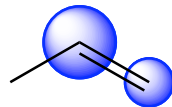
Fixed (VB)

Optimised

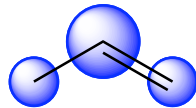
Atomic



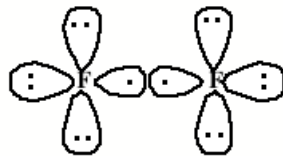
Partial Delocal



Delocal



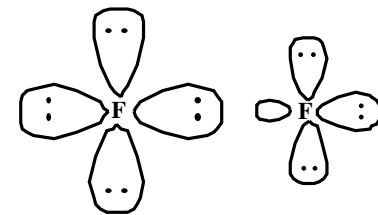
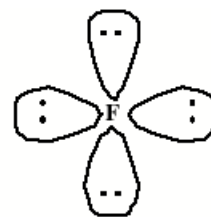
Breathing



covalent



ionic A



ionic B

(overlaps to 99 %)



**G. Balint-Kurti, G. Gallup, R. McWeeny, J. van Lenthe,
M. Raimondi, Y. Mo, W.Wu, R. Harcourt, D. Cooper, J.Gerratt,
R. Braam, , P-A Malmquist, R. Havenith,
S. Shaik, P. Hiberty, S. Humbel, P. Karadov**

+

and a lot of Physics (superconductivity, crystals)

e.g.

Masafumi Tamura, Akiko Nakao, Reizo Kato,

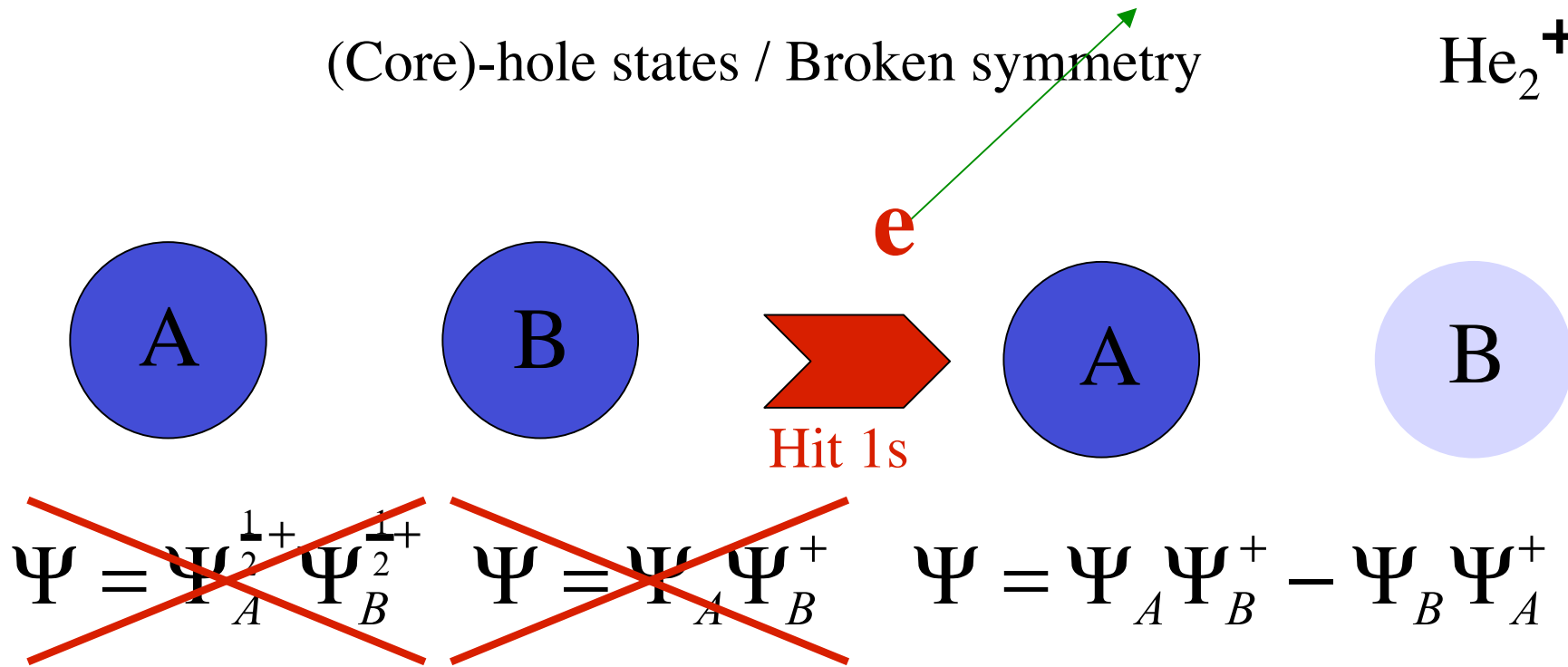
“Frustration-induced valence - bond ordering

in a new quantum triangular antiferromagnet based on [Pd(dmit)₂]. “

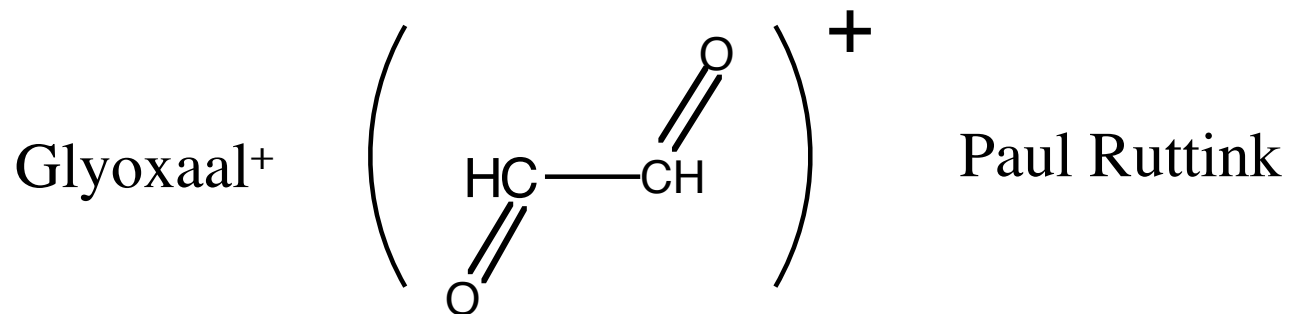
XMVB, CASVB, TURTLE, VB2000, CRUNCH,...

(Core)-hole states / Broken symmetry

He₂⁺



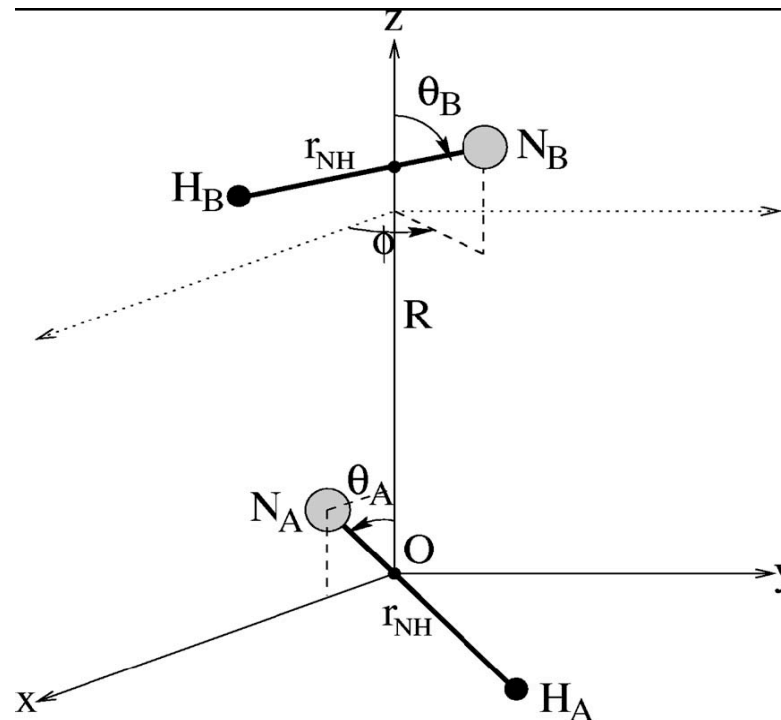
$\langle \Psi_A \Psi_B^+ | \Psi_A^+ \Psi_B \rangle > 0.9$
 Ria Broer - GNOME



NH ($^3\Sigma^-$) dimer

G.Dhont, G. Groenenboom, A. vd Avoird

VB allows states to remain “themselves”
as interaction (and overlap) kicks in



Some wise (edited) remarks by
J.N. Murrell, S.F.A. Kettle and J.M. Tedder
The Chemical Bond
(Wiley, 1985)

Valence bond theory relies on chemical intuition for finding suitable valence structures. In some respect, this may seem an advantage in that chemical experience can be directly introduced into a quantum-chemical calculation. In another respect, however, it is a weakness because it assumes some knowledge of the answer at the start. If chemical intuition is wrong, the VB calculation is wrong. The most valuable calculations are often those whose results refute intuition. In contrast, molecular orbital theory does not make any assumption about the way orbitals are paired in a molecule.

Probably the most misunderstood concept of the theory is “resonance energy”. From resonance theory we deduce that a resonance hybrid will have a lower energy than any single canonical form. The theory tells us nothing about the absolute stability of one species relative to another.